

On eigen-forms on the 2-sphere

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ABSTRACT

We prove that, on the unit 2-sphere, the only eigen-value of the Laplacian on 1-forms less than 4 is equal to 2. The corresponding eigen-forms are described completely.

INTRODUCTION

Let (M, g) be an orientable compact connected n -dimensional Riemannian manifold. To each p -form ω on M , there is associated the $(p+1)$ -form $d\omega$ and the $(n-p)$ -form $*\omega$ respectively, $*$ being the Hodge operator.

The exterior codifferential δ is then defined by

$$\delta\omega = (-1)^p *^{-1} d*\omega, \quad (1)$$

$*^{-1}$ being the inverse mapping to $*$; see Berger *et al.* (1971) and De Rham (1960). Finally, the Laplacian Δ on p -forms is given by

$$\Delta\omega = (d\delta + \delta d)\omega. \quad (2)$$

We say that $\lambda \in \mathbf{R}$ belongs to $\text{Spec}^{(p)}(\Delta)$ if there is a non-trivial p -form ω on M such that

$$\Delta\omega = \lambda\omega. \quad (3)$$

The general problem is to exhibit $\text{Spec}^{(p)}(\Delta)$ for a given (M, g) . Up to now, little is known. $\text{Spec}^{(0)}(\Delta)$ is known just for the hypersphere and several simple homogeneous manifolds. On the other hand, for $n = 2$, $\text{Spec}^{(0)}(\Delta)$ does not contain the intervals

$$(-\infty, 0), (0, 2\kappa_0), (2\kappa_1, 6\kappa_0) \quad (4)$$

with

$$\kappa_0 = \min_M K, \quad \kappa_1 = \max_M K, \quad (5)$$

K being the Gauss curvature of (M, g) . Furthermore, there are no general methods. In what follows, I am going to use just the Stokes theorem

$$\int_M d\phi = 0, \quad (6)$$

ϕ being an $(n-1)$ -form.

LINEAR FORMS ON THE UNIT 2-SPHERE

The purpose of this note is to prove the following:

Theorem. Let (M, g) be the unit sphere of the Euclidean space E^3 , g being the induced metric. Let $\lambda < 4$ and $\lambda \in \text{Spec}^{(1)}(\Delta)$. Then $\lambda = 2$. In E^3 , take two vectors \mathbf{v}, \mathbf{w} . The 1-forms $\omega_{\mathbf{v}}, \omega_{\mathbf{w}}$ on M are defined by

$$\omega_{\mathbf{v}}(t) = \langle \mathbf{v}, t \rangle, \quad \omega_{\mathbf{w}}(t) = \langle \mathbf{w}, t \rangle; \quad t \in T(M); \quad (7)$$

$\langle \cdot, \cdot \rangle$ being the scalar product in E^3 . Then the 1-form

$$\omega = \omega_{\mathbf{v}} + *\omega_{\mathbf{w}} \quad (8)$$

is the most general eigen-form satisfying $\Delta\omega = 2\omega$.

Proof. Let $M \subset E^3$ be the unit sphere; we are going to investigate its coordinate neighbourhood $U \subset M$. To each point $m \in M$, let us associate an orthonormal frame $\{\mathbf{v}_1, \mathbf{v}_2, \mathbf{v}_3\}$ such that $m + \mathbf{v}_3$ is the centre of our sphere. Then we have

$$\begin{aligned} dm &= \omega^1 \mathbf{v}_1 + \omega^2 \mathbf{v}_2, & d\mathbf{v}_1 &= \omega_1^2 \mathbf{v}_2 + \omega^1 \mathbf{v}_3, \\ d\mathbf{v}_2 &= -\omega_1^2 \mathbf{v}_1 + \omega^2 \mathbf{v}_3, & d\mathbf{v}_3 &= -\omega^1 \mathbf{v}_1 - \omega^2 \mathbf{v}_2, \end{aligned} \quad (9)$$

the 1-forms $\omega^1, \omega^2, \omega_1^2$ satisfying

$$d\omega^1 = -\omega^2 \wedge \omega_1^2, \quad d\omega^2 = \omega^1 \wedge \omega_1^2, \quad d\omega_1^2 = -\omega^1 \wedge \omega^2. \quad (10)$$

On M , be given a 1-form ω ; in U , we may write

$$\omega = a\omega^1 + b\omega^2. \quad (11)$$

$a, b: U \rightarrow \mathbf{R}$ being functions. The functions $a_i, b_i: U \rightarrow \mathbf{R}$ are defined by

$$\begin{aligned} da - b\omega_1^2 &= a_1\omega^1 + a_2\omega^2, \\ db + a\omega_1^2 &= b_1\omega^1 + b_2\omega^2. \end{aligned} \quad (12)$$

The exterior differentiation implies

$$\begin{aligned} \{da_1 - (a_2 + b_1)\omega_1^2\} \wedge \omega^1 + \{da_2 + (a_1 - b_2)\omega_1^2\} \wedge \omega^2 &= b\omega^1 \wedge \omega^2, \\ \{db_1 + (a_1 - b_2)\omega_1^2\} \wedge \omega^1 + \{db_2 + (a_2 + b_1)\omega_1^2\} \wedge \omega^2 &= -a\omega^1 \wedge \omega^2 \end{aligned} \quad (13)$$

and the existence of functions $a_{ij}, b_{ij}: U \rightarrow \mathbf{R}$ such that

$$\begin{aligned} da_1 - (a_2 + b_1)\omega_1^2 &= a_{11}\omega^1 + (a_{12} - \frac{1}{2}b)\omega^2, \\ da_2 + (a_1 - b_2)\omega_1^2 &= (a_{12} + \frac{1}{2}b)\omega^1 + a_{22}\omega^2, \\ db_1 + (a_1 - b_2)\omega_1^2 &= b_{11}\omega^1 + (b_{12} + \frac{1}{2}a)\omega^2, \\ db_2 + (a_2 + b_1)\omega_1^2 &= (b_{12} - \frac{1}{2}a)\omega^1 + b_{22}\omega^2. \end{aligned} \quad (14)$$

The differential consequences of (14) being

$$\begin{aligned} \{da_{11} - (2a_{12} + b_{11})\omega_1^2\} \wedge \omega^1 + \{da_{12} + (a_{11} - a_{22} - b_{12})\omega_1^2\} \wedge \omega^2 &= (a_2 + \frac{3}{2}b_1)\omega^1 \wedge \omega^2, \\ \{da_{12} + (a_{11} - a_{22} - b_{12})\omega_1^2\} \wedge \omega^1 + \{da_{22} + (2a_{12} - b_{22})\omega_1^2\} \wedge \omega^2 &= (\frac{3}{2}b_2 - a_1)\omega^1 \wedge \omega^2, \\ \{db_{11} + (a_{11} - 2b_{12})\omega_1^2\} \wedge \omega^1 + \{db_{12} + (a_{12} + b_{11} - b_{22})\omega_1^2\} \wedge \omega^2 &= (b_2 - \frac{3}{2}a_1)\omega^1 \wedge \omega^2, \\ \{db_{12} + (a_{12} + b_{11} - b_{22})\omega_1^2\} \wedge \omega^1 + \{db_{22} + (a_{22} + 2b_{12})\omega_1^2\} \wedge \omega^2 &= -(\frac{3}{2}a_2 + b_1)\omega^1 \wedge \omega^2. \end{aligned} \quad (15)$$

we get the existence of functions $A_x, B_x: U \rightarrow \mathbf{R}$ satisfying

$$\begin{aligned}
da_{11} - (2a_{12} + b_{11})\omega_1^2 &= A_1\omega^1 + (A_2 - \frac{1}{2}a_2 - \frac{3}{4}b_1)\omega^2, \\
da_{12} + (a_{11} - a_{22} - b_{12})\omega_1^2 &= (A_2 + \frac{1}{2}a_2 + \frac{3}{4}b_1)\omega^1 + (A_3 + \frac{1}{2}a_1 - \frac{3}{4}b_2)\omega^2, \\
da_{22} + (2a_{12} - b_{22})\omega_1^2 &= (A_3 - \frac{1}{2}a_1 + \frac{3}{4}b_2)\omega^1 + A_4\omega^2, \\
db_{11} + (a_{11} - 2b_{12})\omega_1^2 &= B_1\omega^1 + (B_2 + \frac{3}{4}a_1 - \frac{1}{2}b_2)\omega^2, \\
db_{12} + (a_{12} + b_{11} - b_{22})\omega_1^2 &= (B_2 - \frac{3}{4}a_1 + \frac{1}{2}b_2)\omega^1 + (B_3 + \frac{3}{4}a_2 + \frac{1}{2}b_1)\omega^2, \\
db_{22} + (a_{22} + 2b_{12})\omega_1^2 &= (B_3 - \frac{3}{4}a_2 - \frac{1}{2}b_1)\omega^1 + B_4\omega^2.
\end{aligned} \tag{16}$$

For 1-forms, we have

$$*(p\omega^1 + q\omega^2) = -q\omega^1 + p\omega^2, \quad *^{-1}(p'\omega^1 + q'\omega^2) = q'\omega^1 - p'\omega^2; \tag{17}$$

$$\Delta\omega = (*^{-1}d*d - d*^{-1}d*)\omega. \tag{18}$$

In our case,

$$\begin{aligned}
d\omega &= (b_1 - a_2)\omega^1 \wedge \omega^2, & *d\omega &= b_1 - a_2, \\
d*d\omega &= (b_{11} - a_{12} - \frac{1}{2}b)\omega^1 + (b_{12} - a_{22} + \frac{1}{2}a)\omega^2, \\
*^{-1}d*d\omega &= (b_{12} - a_{22} + \frac{1}{2}a)\omega^1 + (a_{12} - b_{11} + \frac{1}{2}b)\omega^2, \\
\omega &= -b\omega^1 + a\omega^2, & d\omega &= (a_1 + b_2)\omega^1 \wedge \omega^2, & *^{-1}d*\omega &= a_1 + b_2, \\
d*^{-1}d*\omega &= (a_{11} + b_{12} - \frac{1}{2}a)\omega^1 + (a_{12} + b_{22} - \frac{1}{2}b)\omega^2.
\end{aligned} \tag{19}$$

Thus, for the 1-form ω (11),

$$\Delta\omega = (a - a_{11} - a_{22})\omega^1 + (b - b_{11} - b_{22})\omega^2. \tag{20}$$

The form ω satisfying (6), we get

$$a_{11} + a_{22} = (1 - \lambda)a, \quad b_{11} + b_{22} = (1 - \lambda)b. \tag{21}$$

Because of (16), we obtain

$$\begin{aligned}
A_1 + A_3 &= (\frac{3}{2} - \lambda)a_1 - \frac{3}{4}b_2, & A_2 + A_4 &= (\frac{3}{2} - \lambda)a_2 + \frac{3}{4}b_1, \\
B_1 + B_3 &= \frac{3}{4}a_2 + (\frac{3}{2} - \lambda)b_1, & B_2 + B_4 &= -\frac{3}{4}a_1 + (\frac{3}{2} - \lambda)b_2.
\end{aligned} \tag{22}$$

For a general 1-form ω (11), we easily get

$$\begin{aligned}
d*d\{(a_1 - b_2)^2 + (a_2 + b_1)^2\} &= \\
&= [2\{(a_{11} - b_{12} + \frac{1}{2}a)^2 + (a_{12} + b_{11} + \frac{1}{2}b)^2 + (a_{12} - b_{22} - \frac{1}{2}b)^2 + (a_{22} + b_{12} + \frac{1}{2}a)^2\} + \\
&\quad + \frac{7}{2}\{(a_1 - b_2)^2 + (a_2 + b_1)^2\} + \\
&\quad + 2\{(a_1 - b_2)(A_1 + A_3 - B_2 - B_4) + (a_2 + b_1)(A_2 + A_4 + B_1 + B_3)\}]\omega^1 \wedge \omega^2.
\end{aligned} \tag{23}$$

For an eigen-form, we have (22) and

$$\begin{aligned}
d*d\{(a_1 - b_2)^2 + (a_2 + b_1)^2\} &= \\
&= [2\{(a_{11} - b_{12} + \frac{1}{2}a)^2 + (a_{12} + b_{11} + \frac{1}{2}b)^2 + (a_{12} - b_{22} - \frac{1}{2}b)^2 + (a_{22} + b_{12} + \frac{1}{2}a)^2\} + \\
&\quad + 2(4 - (a_1 - b_2)^2 + (a_2 + b_1)^2)]\omega^1 \wedge \omega^2.
\end{aligned} \tag{24}$$

Using the Stokes theorem and the supposition $4 - \lambda > 0$, we get

$$a_1 - b_2 = 0, \quad a_2 + b_1 = 0, \quad (25)$$

$$a_{11} - b_{12} + \frac{1}{2}a = a_{12} + b_{11} + \frac{1}{2}b = a_{12} - b_{22} - \frac{1}{2}b = a_{22} + b_{12} + \frac{1}{2}a = 0; \quad (26)$$

It is easy to see that (26) is just a differential consequence of (25).

Let us introduce the functions $r, s: U \rightarrow \mathbf{R}$ by

$$r = a_1 = b_2, \quad s = -a_2 = b_1; \quad (27)$$

the equation (12) is then

$$\begin{aligned} da - b\omega_1^2 &= r\omega^1 - s\omega^2, \\ db + a\omega_1^2 &= s\omega^1 + r\omega^2. \end{aligned} \quad (28)$$

Using successive differentiations and the Cartan's lemma, we get the existence of functions $x, y, t, z, p, q: U \rightarrow \mathbf{R}$ such that

$$\begin{aligned} dr \wedge \omega^1 - ds \wedge \omega^2 &= b\omega^1 \wedge \omega^2, \\ ds \wedge \omega^1 + dr \wedge \omega^2 &= -a\omega^1 \wedge \omega^2, \end{aligned} \quad (29)$$

$$\begin{aligned} dr &= (x - \frac{1}{2}a)\omega^1 - (y + \frac{1}{2}b)\omega^2, \\ ds &= (y - \frac{1}{2}b)\omega^1 + (x + \frac{1}{2}a)\omega^2; \end{aligned} \quad (30)$$

$$\begin{aligned} (dx + y\omega_1^2) \wedge \omega^1 - (dy - x\omega_1^2) \wedge \omega^2 &= s\omega^1 \wedge \omega^2, \\ (dy - x\omega_1^2) \wedge \omega^1 + (dx + y\omega_1^2) \wedge \omega^2 &= -r\omega^1 \wedge \omega^2, \end{aligned} \quad (31)$$

$$\begin{aligned} dx + y\omega_1^2 &= (t - \frac{1}{2}r)\omega^1 - (z + \frac{1}{2}s)\omega^2, \\ dy - x\omega_1^2 &= (z - \frac{1}{2}s)\omega^1 + (t + \frac{1}{2}r)\omega^2; \end{aligned} \quad (32)$$

$$\begin{aligned} (dt + 2z\omega_1^2) \wedge \omega^1 - (dz - 2t\omega_1^2) \wedge \omega^2 &= 0, \\ (dz - 2t\omega_1^2) \wedge \omega^1 + (dt + 2z\omega_1^2) \wedge \omega^2 &= 0, \end{aligned} \quad (33)$$

$$\begin{aligned} dt + 2z\omega_1^2 &= p\omega^1 - q\omega^2, \\ dz - 2t\omega_1^2 &= q\omega^1 + p\omega^2. \end{aligned} \quad (34)$$

Now,

$$d * d \{ (x - \frac{1}{2}a)^2 + (y + \frac{1}{2}b)^2 + r^2 \} = 4(z^2 + t^2)\omega^1 \wedge \omega^2, \quad (35)$$

and the Stokes theorem implies

$$z = t = 0. \quad (36)$$

Let us consider the vectors

$$\begin{aligned} \mathbf{v} &= (\frac{1}{2}a - x)\mathbf{v}_1 + (y + \frac{1}{2}b)\mathbf{v}_2 + r\mathbf{v}_3, \\ \mathbf{w} &= (\frac{1}{2}b - y)\mathbf{v}_1 - (x + \frac{1}{2}a)\mathbf{v}_2 + s\mathbf{v}_3; \end{aligned} \quad (37)$$

It is easy to see that

$$d\mathbf{v} = d\mathbf{w} = 0, \quad (38)$$

i.e., \mathbf{v} and \mathbf{w} are fixed vectors of E^3 . The forms (7) are then

$$\omega_{\mathbf{v}} = (\frac{1}{2}a - x)\omega^1 + (y + \frac{1}{2}b)\omega_1^2, \quad \omega_{\mathbf{w}} = (\frac{1}{2}b - y)\omega^1 - (x + \frac{1}{2}a)\omega^2, \quad (39)$$

and we get exactly (8).

Finally, let

$$\mathbf{v} = e\mathbf{v}_1 + f\mathbf{v}_2 + g\mathbf{v}_3, \quad \mathbf{w} = k\mathbf{v}_1 + l\mathbf{v}_2 + m\mathbf{v}_3 \quad (40)$$

be fixed vectors in E^3 . From

$$0 = d\mathbf{v} = (de - f\omega_1^2 - g\omega^1)\mathbf{v}_1 + (df + e\omega_1^2 - g\omega^2)\mathbf{v}_2 + (dg + e\omega^1 + f\omega^2)\mathbf{v}_3, \quad (41)$$

we get

$$de - f\omega_1^2 = g\omega^1, \quad df + e\omega_1^2 = g\omega^2, \quad dg = -e\omega^1 - f\omega^2 \quad (42)$$

and similar equations

$$dk - l\omega_1^2 = m\omega^1, \quad dl + k\omega_1^2 = m\omega^2, \quad dm = -k\omega^1 - l\omega^2 \quad (43)$$

for \mathbf{w} . We have

$$\omega_{\mathbf{v}} = e\omega^1 + f\omega^2, \quad \omega_{\mathbf{w}} = k\omega^1 + l\omega^2, \quad (44)$$

i.e.

$$\omega = \omega_{\mathbf{v}} + * \omega_{\mathbf{w}} = (e + l)\omega^1 + (f + k)\omega^2; \quad (45)$$

$$d\omega = 2m\omega^1 \wedge \omega^2, \quad *d\omega = 2m, \quad d*d\omega = -2k\omega^1 - 2l\omega^2,$$

$$*^{-1}d*d\omega = -2l\omega^1 + 2k\omega^2;$$

$$*\omega = -(f + k)\omega^1 + (e - l)\omega^2, \quad d*\omega = 2g\omega^1 \wedge \omega^2, \quad *^{-1}d*\omega = 2g,$$

$$d*^{-1}d*\omega = -2e\omega^1 - 2f\omega^2, \quad (46)$$

i.e.

$$\Delta\omega = 2(e - l)\omega^1 + 2(f + k)\omega^2 = 2\omega. \quad (47)$$

This proves our theorem.

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